



## Erratum: Strong coupling between weakly guided semiconductor nanowire modes and an organic dye [Phys. Rev. B **99**, 205409 (2019)]

Diego R. Abujetas , Johannes Feist, Francisco J. García-Vidal, Jaime Gómez Rivas, and José A. Sánchez-Gil

 (Received 26 November 2020; published 18 December 2020)

DOI: [10.1103/PhysRevB.102.239901](https://doi.org/10.1103/PhysRevB.102.239901)

There were errors in the reproduction of several equations in Sec. III of the original paper. None of the conclusions or figures are affected as they were obtained using the correct equations given in the following.

Equation (5) should read

$$U_{\text{rad}} = 2\varepsilon_0 \int \varepsilon(r)r|\vec{E}(r)|^2 dr. \quad (5)$$

Equations (7) and (8) and the accompanying text should be as follows:

$$H_{\text{mol},j} = \hbar \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & \omega_1 - i\frac{\gamma_1}{2} & 0 & 0 \\ 0 & 0 & \omega_2 - i\frac{\gamma_2}{2} & 0 \\ 0 & 0 & 0 & \omega_3 - i\frac{\gamma_3}{2} \end{pmatrix}, \quad (7)$$

$$\hat{\mu}_j = \begin{pmatrix} 0 & \mu_1 & \mu_2 & \mu_3 \\ \mu_1 & 0 & 0 & 0 \\ \mu_2 & 0 & 0 & 0 \\ \mu_3 & 0 & 0 & 0 \end{pmatrix}, \quad (8)$$

where the parameters  $\omega_k$  and  $\gamma_k$  are taken from the fit in Eq. (3). The transition dipole moments  $\mu_k$  can be related to the amplitude parameters  $a_k$  in Eq. (3) through  $\mu_k^2 = 3\varepsilon_0\hbar a_k/\rho_{\text{mol}}$ , where  $\rho_{\text{mol}}$  is the number density of molecules, by inserting the polarizability of the model molecules into the Clausius-Mossotti relation while neglecting dipole-dipole interactions (i.e., setting  $\varepsilon + 2 \approx 3$ ). Consequently, we do not include direct dipole-dipole interactions between the molecules, as their (averaged) effect is already included in the parameters extracted from the permittivity.

Finally, Eqs. (9)–(13) should be as follows:

$$g_{nmk}^2 = \sum_j |\vec{\mathcal{E}}_{nm}(\vec{r}_j) \cdot \vec{d}_j \mu_k|^2. \quad (9)$$

$$g_{nmk}^2 \approx \frac{2\pi L}{3} \frac{\mu_k^2}{V_{\text{mol}}} \int_R^\infty r |\vec{\mathcal{E}}_{nm}(r)|^2 dr. \quad (10)$$

$$g_{nmk}^2 = \frac{\hbar\omega_{c,nm}}{3\varepsilon_0\varepsilon_h} \frac{\mu_k^2}{V_{\text{mol}}} \frac{U_{\text{ext},nm}^{el}}{U_{\text{rad},nm}}, \quad (11)$$

$$U_{\text{ext}}^{el} = \varepsilon_0\varepsilon_h \int_R^\infty r |\vec{E}(r)|^2 dr. \quad (12)$$

$$g_{nmk} = \sqrt{\frac{\mu_k^2 \rho_{\text{mol}}}{3\varepsilon_0\varepsilon_h} \hbar\omega_{c,nm} F_{\text{ext},nm}} = \sqrt{\frac{\hbar a_k}{\varepsilon_h} \hbar\omega_{c,nm} F_{\text{ext},nm}}. \quad (13)$$